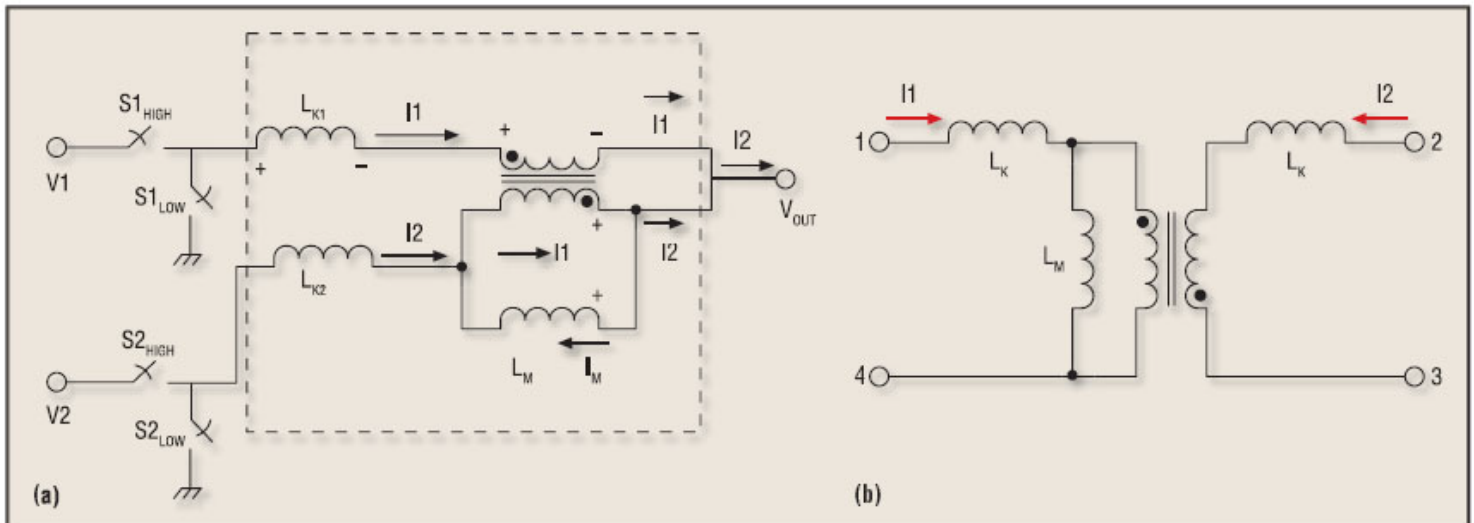


## DESIGNING COUPLED INDUCTORS

### Power Electronics

Using a previously derived circuit model, coupled inductor designs can be optimized for best performance in multiphase power converters.

In the January issue of Power Electronics Technology, the operation and benefits of the coupled inductor multiphase (CIMP) topology were reviewed, and the equations for the output and phase-ripple currents were derived using a circuit model representation of the coupled inductor. In the circuit model, the coupled inductor is represented by a leakage inductance in each phase ( $L_K$ ), a magnetizing inductance ( $L_M$ ) and an ideal transformer with a 1-to-1 turns ratio (Fig. 1).



**Fig. 1.** A simplified schematic of the coupled inductor multiphase (CIMP) topology is shown (a). The coupled inductor schematic (b) is identical to that of any two-winding transformer. The equivalent resistive losses (due to core loss and copper loss) as well as interwinding capacitances are ignored for simplicity. The magnetizing inductance ( $L_M$ ) can be placed on either the primary or secondary winding with no effect to the circuit equations due to the 1-to-1 turns ratio.

It was determined that the principle benefit of the CIMP topology was a significant reduction in phase-ripple current for a given transient response condition when compared to the uncoupled version of this topology. It was shown that this ripple reduction was dependant on the duty cycle ( $D$ ), which is  $V_{OUT}/V_{IN}$ , and the coupling ratio, which is  $p = L_M/L_K$ . This reduction allows for improved efficiency or, conversely, one can improve the transient response using the CIMP topology and still maintain the same efficiency as in the uncoupled variety. It is clear that the CIMP topology relies on the design and implementation of the coupled inductor component. By leveraging the previous analysis, we can determine the relevant equations required to design and optimize the coupled inductor component.

## The Coupled Inductor

Any inductor or transformer with multiple windings is a coupled inductor. Therefore, hundreds of implementations exist. However, in the vast majority of these implementations, with the notable exception of the common-mode choke, the driving source is only applied to the primary — or main — winding, and the other windings are simply “dumb” followers mimicking the behavior of the first.

Although, in theory, these devices could be used in a CIMP application, they would not work well. The trick is to determine how to design and optimize a part that has the required magnetizing and leakage inductances, appropriate saturation and thermal characteristics in a package size that is suitable to the application. To do this, it is necessary to first identify some possible core structures and their corresponding reluctance models, and then use the reluctance model in conjunction with the coupled inductor circuit model to realize a complete solution.

Although there is an infinite number of possible core structures for a coupled inductor, the toroid and E-core designs are two that immediately come to mind as shown in Figs. 2a and 2b. A reluctance model is the magnetic equivalent of an electric circuit model in which a magnetic force (analogous to voltage) drives a magnetic field or flux (analogous to current) through a reluctance path (analogous to resistance), and as such any electrical theorems or rules apply similarly.

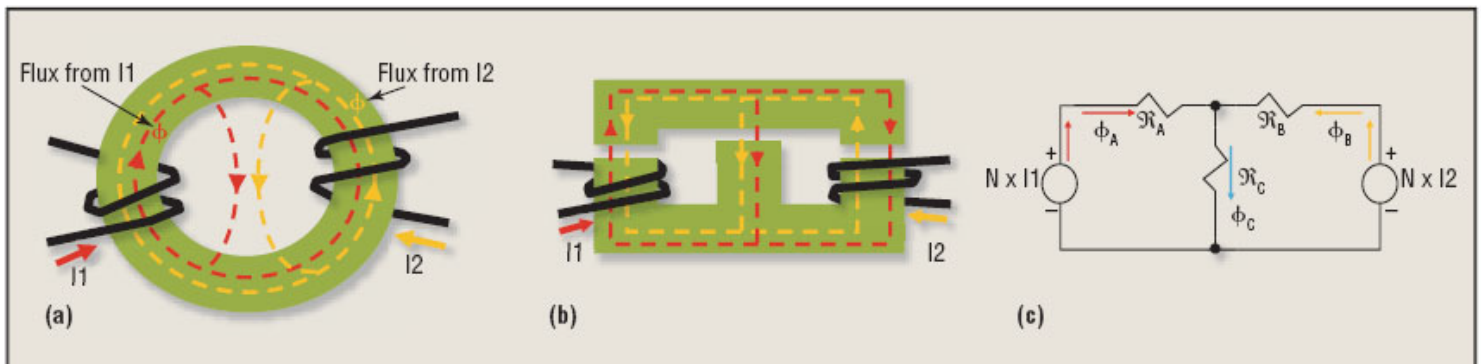


Fig. 2. Both the toroid design (a) and the E-core design (b) can be reduced to the same reluctance model (c).

The easiest way to determine a reluctance model for a coupled inductor is to remove one of the windings and then envision the various paths that magnetic flux could take to complete a closed loop back to the driving force (magnetic force from the remaining winding is equal to the number of turns times the driving current). Each flux path has an equivalent reluctance that is equal to the length of the flux path ( $l_E$ ) divided by the cross-sectional area of the path ( $A_E$ ) and the permeability of the material ( $\mu$ ), or  $R = l_E / (A_E \times \mu)$ . The lower the reluctance for a given path, the greater the magnetic flux through that path will be.

Although it is possible to develop a model in which every flux path is identified, it is easier and just as accurate to identify the most likely paths and use these in the reluctance model. It is then possible to refine the model by adding the less likely paths back in, if desired. After identifying the most likely flux paths, the other winding is re-inserted into the reluctance model and the model is complete.

As can be seen in Fig. 2a, most of the flux  $\Phi_1$  (red) created by the magnetic force,  $N \times I_1$ , travels through the high-permeability core (low reluctance), although there is some leakage flux that takes the shorter but higher reluctance path through the air. Although the flux leaks everywhere, all of these reluctance paths are in parallel and can be equated with a single path through the center of the core. The same effect happens with the flux  $\Phi_2$  (yellow) from the  $N \times I_2$  magnetic source. In Fig. 2b, intentional air gaps are inserted into the high-permeability core material to help direct the flux. Although there is flux that travels outside the core geometry, it is comparatively small and can be ignored.

As shown in Fig. 2c, although the structures in Figs. 2a and 2b are physically different, the reluctance model for each is the same. It can be assumed that the construction of the inductor is symmetrical and, therefore, the reluctances  $R_A$  and  $R_B$  are equal and can be replaced with  $R$ . For each of the identified structures, the values of reluctances  $R$  and  $R_C$  will vary depending on the placement of the windings, the core material used, the size of any inserted air gaps and the physical size of the core. The trick now is to determine how reluctances  $R$  and  $R_C$  from the reluctance model relate to the magnetizing inductance and the series leakage inductances  $L_M$  and  $L_K$  of the circuit model.

One approach to relate reluctances  $R$  and  $R_c$  to the circuit elements  $L_M$  and  $L_K$  is to look at how one would actually measure the circuit elements (Fig. 3a), and how this measurement equates back to the reluctance model. The first and most obvious measurement to take on the inductor component is to measure the open-circuit inductance on each of the two windings. Fig. 3b makes it clear that the measured open-circuit inductance equals:

$$(Eq. 1) \quad L_{OPEN(1-4)} = L_{OPEN(2-3)} = L_M + L_K$$

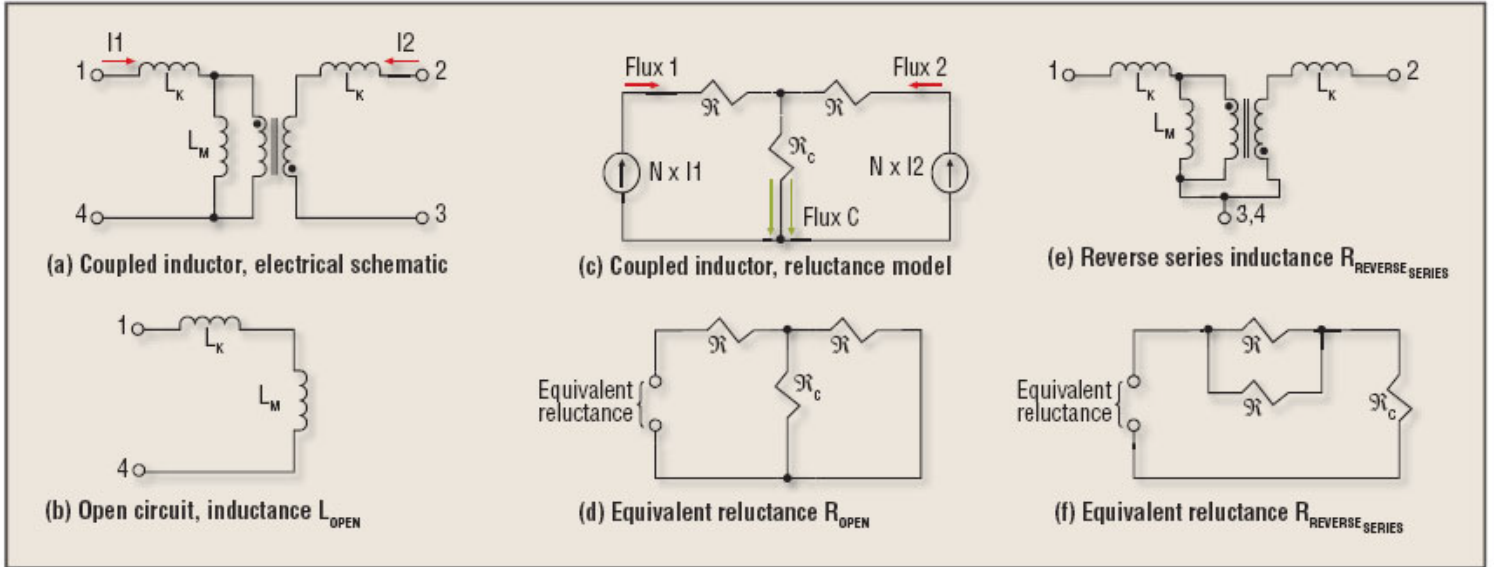


Fig.3. In order to write the equations for the circuit elements  $L_M$  and  $L_K$  in terms of the reluctance elements  $R$  and  $R_c$ , it is helpful to analyze several possible inductance measurements and how these measurements affect the reluctance model.

In regards to the reluctance model (Fig. 3c), the open circuit inductance measurement on only one winding is equivalent to removing or shorting the other source such that the equivalent circuit reduces to that shown in Fig.3d. It is clear from this figure that the equivalent reluctance is equal to:

$$R_{OPEN(1-4)} = R_{OPEN(2-3)} =$$

$$R + (R_c \parallel R) = \frac{(2RR_c + R^2)}{R_c + R}$$

( Eq. 2 )

Knowing that  $L = \frac{N^2}{R}$  and using Eqs. 1 and 2 yields:

$$L_{OPEN(1-4)} = L_{OPEN(2-3)} = L_M + L_K = N^2 \frac{(R + R_c)}{(2RR_c + R^2)}$$

( Eq. 3 )

The second measurement on the inductor component could be to measure the short-circuit inductance by measuring the inductance on one winding with the other winding shorted. Assuming a perfect short, the short-circuit inductance equals:

$$L_{SHORT(1-4)} \text{ with (2-3) shorted} = L_K + (L_M \parallel L_K),$$

( Eq. 4 )

$$\text{If } L_M \gg L_K, \text{ then } L_{SHORT(1-4)} = 2 \times L_K$$

This measurement is often used to measure the leakage inductance ( $L_K$ ) in transformer applications because it is assumed that  $L_M \gg L_K$ . However, for the coupled inductor, this is not the case, and as such  $L_{SHORT}$  does not lead to a clear or direct measurement of the leakage inductance. A better measurement is the reverse-series inductance, in which the windings of the inductor are tied in series but out of phase (Fig. 3e). In this measurement, the opposing polarity of the windings effectively cancels the magnetizing inductance ( $L_M$ ) and the series inductance equals:

$$(Eq. 5) \quad L_{REVERSE\_SERIES(1-2)} \text{ (with 3-4 shorted)} = 2L_K.$$

In regards to the reluctance model, the reverse-series measurement is equivalent to putting the two windings in parallel, and the equivalent circuit reduces to that shown in Fig. 3f. It is clear from this figure that the equivalent reluctance is equal to:

$$(Eq. 6) \quad \mathfrak{R}_{REVERSE\_SERIES(1-2)} \text{ (with 3-4 shorted)} = \mathfrak{R}_C + (\mathfrak{R} \parallel \mathfrak{R}) = \mathfrak{R}_C + 0.5\mathfrak{R},$$

Again, knowing that  $L = N^2 / R$  and using Eqs. 5 and 6 yields:

$$(Eq. 7) \quad L_{REVERSE\_SERIES(1-2)} \text{ (with 3-4 shorted)} = 2 \times L_K = \frac{N^2}{(\mathfrak{R}_C + 0.5\mathfrak{R})},$$

$$(Eq. 8) \quad \text{or } L_K = \frac{N^2}{(\mathfrak{R} + 2\mathfrak{R}_C)}.$$

And using Eq. 8 to solve Eq. 3 yields:

$$(Eq. 9) \quad L_M = N^2 \frac{\mathfrak{R}_C}{(2\mathfrak{R}\mathfrak{R}_C + \mathfrak{R}^2)}.$$

With the definition of  $L_K$  and  $L_M$  in terms of  $R$  and  $R_C$  it would be possible now to design a coupled inductor. However, it would not be possible to determine if the inductor was going to saturate in an actual application without first determining the amount of flux in the various sections of the core. To determine this, we need to solve for the flux ( $\Phi_1$ ,  $\Phi_2$  and  $\Phi_3$ ) in the reluctance model. The reluctance model (Fig. 3c) resembles the electrical circuit used in Millman's Theorem, and as such this theorem can be used to solve for the flux in each leg:

$$(Eq. 10) \quad \Phi_1 = \left( \frac{(\mathfrak{R} + \mathfrak{R}_C)}{(\mathfrak{R}^2 + 2\mathfrak{R}\mathfrak{R}_C)} \right) (N \times I_1) - \frac{\mathfrak{R}_C}{(\mathfrak{R}^2 + 2\mathfrak{R}\mathfrak{R}_C)} (N \times I_2)$$

$$(Eq. 12) \quad \Phi_3 = \Phi_1 = \Phi_2$$

$$(Eq. 11) \quad \Phi_2 = \frac{-\mathfrak{R}_C}{(\mathfrak{R}^2 + 2\mathfrak{R}\mathfrak{R}_C)} (N \times I_1) + \left( \frac{(\mathfrak{R} + \mathfrak{R}_C)}{(\mathfrak{R}^2 + 2\mathfrak{R}\mathfrak{R}_C)} \right) (N \times I_2)$$

where I1 and I2 are the currents in phase one and phase two, respectively.

It can be seen that the expression  $R_c / (R^2 + 2RR_c)$  in Eqs. 10 and 11 is similar to the definition of  $L_M$  in Eq. 9, and it should be noted that the expression  $(R + R_c) / (R^2 + 2RR_c)$  is similar to adding  $L_M$  plus  $L_K$  in Eq. 3. Substitution and noting that  $L_M = p \times L_K$  yields:

$$\begin{aligned} \text{( Eq. 13 )} \quad \phi_1 &= \left( \frac{(L_M + L_K)}{N} \right) \times I1 - \left( \frac{L_M}{N} \right) \times I2 = \\ &\left( \frac{L_K}{N} \right) \times I1 + \left( \frac{pL_K}{N} \right) \times (I1 - I2) \end{aligned}$$

$$\begin{aligned} \text{( Eq. 14 )} \quad \phi_2 &= -\left( \frac{L_M}{N} \right) \times I1 + \left( \frac{(L_M + L_K)}{N} \right) \times I2 = \\ &\left( \frac{L_K}{N} \right) \times I2 + \left( \frac{pL_K}{N} \right) \times (I2 - I1) \end{aligned}$$

$$\text{( Eq. 15 )} \quad \phi_3 = \left( \frac{L_K}{N} \right) \times (I1 + I2) = \left( \frac{L_K}{N} \right) \times I_{OUT}$$

Three things should be noted about these equations. First, the flux in the center of the reluctance model ( $\Phi_3$ ) is not dependent on the coupling factor (p) or the imbalance between the currents in phases one and two (I1 and I2). Second, and most important, the flux in the outer legs of the reluctance model ( $\Phi_1$  and  $\Phi_2$ ) is very much dependent on the current imbalance as well as the coupling factor (p). As a result, high values of p (as p is increased, the phase ripple decreases but at a decreasing rate) actually destabilize the inductor and make it more susceptible to saturation caused by a current imbalance between phases; therefore, high p values should be avoided.

By substituting the phase current equations derived in the January 2006 article, Eqs. 13, 14 and 15 can be rewritten in terms of the peak flux (affecting saturation) and the ac flux (affecting core losses) as:

$$\text{( Eq. 16 )} \quad \phi_{1PK} = \phi_{2PK} = \frac{L_K}{N} I_{1PK} + \frac{pL_K}{N} I_{CURRENT\_IMBALANCE}$$

$$\text{( Eq. 17 )} \quad \phi_{2PK} = \frac{L_K}{N} I_{OUTPK}$$


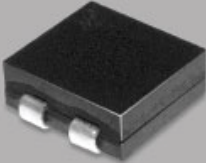
$$\text{( Eq. 18 )} \quad \phi_{1AC} = \phi_{2AC} = \frac{V_{IN}}{N} (1 - D) DT_s$$

$$\text{( Eq. 19 )} \quad \phi_{3AC} = \frac{V_{IN}}{N} (1 - 2D) DT_s$$

Note that the  $\Phi_{3AC}$  occurs at twice the switching frequency ( $F_s$ ), but that the flux in the outer legs  $\Phi_{1AC}$  and  $\Phi_{2AC}$  occurs at the switching frequency. Intuitively, this makes sense because the outer legs are only being driven once per cycle, whereas the center leg is being driven by both the  $N \times I1$  and  $N \times I2$  terms. In addition, note that the ac flux is not affected by the coupling ratio (p) or the actual value of  $L_K$  as it is influenced entirely by the applied volt- $\mu$ s in the application.

## Uncoupled vs Coupled Design

To get a feel for the tradeoffs between uncoupled and coupled inductor applications, it is helpful to run through an actual design comparison. A potential application for the CIMP topology would be in powering the processor for IMVP6 notebook computing applications. Although the electrical requirements vary, a standard specification is shown in the table.

Independent Inductors	IMVP6 Notebook Specification	Coupled Inductors
	Topology: Two-Phase Buck $V_{IN} = 12\text{ V to }19\text{ V}$ Frequency = 300 kHz $V_{OUT} = 1.2\text{ V}$ $I_{OUT} = 44\text{ A dc}$	
PG0255.401NL	Part Number	Pulse
2	Quantity/Board	1
360	Inductance (nH) - $L_K$	676
N/A	Inductance (nH) - $L_M$	183
1.0	DCR (m $\Omega$ )	1.0
11.5 x 10.3	Footprint (mm max)	18.4 x 8.8
4.0	Height (mm max)	4.0
10	Phase Ripple (A pk-pk)	10
30	di/dt max (A/ $\mu$ s)	60 (2x increase)
237	Total Board Area (mm <sup>2</sup> )	162 (32% reduction)
2.0 (2 x 1.0)	Total Relative Cost	1.3 (1 x 1.3) (35% decrease)
1.7 (2 x 0.85)	Power Loss in Core (mW)	0.063

**Table.** Two-phase coupled inductor design versus uncoupled two-inductor solution.

Existing implementations use inductors similar to the Pulse PG0255.401NL (360 nH, 1 m $\Omega$ , 11.2 mm  $\times$  10 mm  $\times$  4 mm), and the electrical and mechanical specifications are shown in the table. For the sake of comparison, the design will seek to keep the same efficiency per phase (equal  $\Delta I_{PHASE}$ ), but improve the transient response time by a factor of two ( $di/dt_{MAXIMUM} = 60\text{ A}/\mu\text{s}$ ). This requires an  $L_K$  of  $0.5 \times 360\text{ nH}$  or 180 nH. For consistency, we will assume that the 1 m is a necessary condition due to the use of inductor current sensing — sensing the current by measuring the voltage across the inductor resistance — and that the 4-mm maximum height is required for this notebook application.

# Designing Coupled Inductors



Note that a patent on the coupled inductor multiphase topology has been granted to Volterra, and so it is my understanding that at the present time the coupled inductor multiphase topology can only be used with approval by or license from Volterra. The coupled inductor product defined previously was developed as a comparison and is not commercially available.

As can be seen in the table, the coupled inductor designed using the previous equations enables a 2x increase in the transient response without affecting the ripple current per phase. This performance improvement is accompanied by a reduction in total inductor footprint of more than 30% and a cost reduction of more than 35%. In addition, because the coupled inductor is made using a ferrite-core material, the core losses can be reduced by a total of 1.7 W, which represents a 2% to 3% increase in system efficiency.

Through simple analysis of the circuit and reluctance models, design equations for a two-phase coupled inductor have been derived. This analysis can be extended to higher phase counts. The proposed two-phase coupled inductor design for notebook applications showed that the CIMP topology can enable faster transient response, improved efficiency, smaller board space and reduced cost over the conventional two inductor noncoupled multiphase approach.

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